# Task space controller for the novel Active Ankle

### Introduction

ACTIVE ANKLE is a novel parallel manipulator with three degrees of freedom that operates in an almost-spherical manner [1, 2]. The almost-spherical parallel manipulator (ASPM) is primarily intended as an actuated ankle joint in a full-body exoskeleton for rehabilitation application (Fig. 3).

### Design features

- 1. lightweight and robust construction
- 2. modular design leading to low link diversity
- 3. high stiffness and orientation accuracy
- 4. high payload capacity
- 5. no torques required for loads along torsional axis



Fig.1: ACTIVE ANKLE prototype

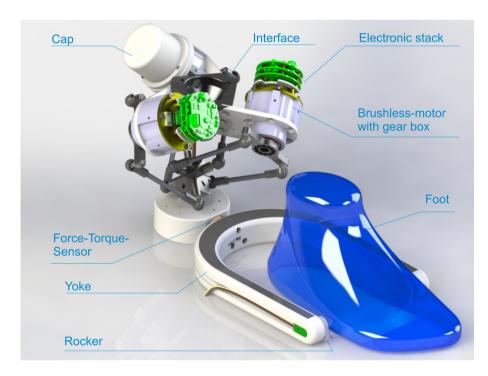


Fig.3: ACTIVE ANKLE with foot unit

Fig.2: Sketch of the ACTIVE ANKLE

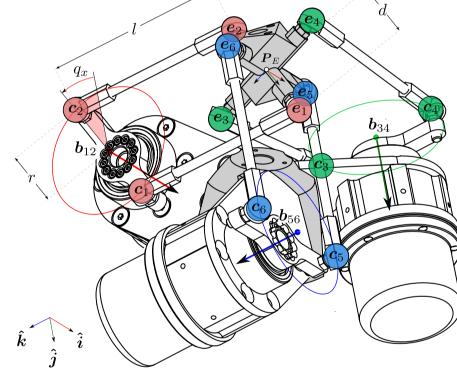


Fig.4: Scheme, r = d = 35, l = 100.

### Control challenge

Due to spatial behaviour but spherical use case of the ACTIVE ANKLE, the task space control of this mechanism asks for a joint configuration for a given orientation from SO(3), instead of a pose from SE(3) [3].

### Inverse Geometric Model (IGM)

The Inverse Geometric Model (IGM) is a solution to the problem of finding input joint angles  $[q_x, q_y, q_z]$  for a specific

end-effector pose  $\mathbf{P}_E = \begin{bmatrix} \mathbf{s} \ \mathbf{n} \ \mathbf{a} \ \mathbf{e} \\ 0 \ 0 \ 1 \end{bmatrix} \in SE(3)$ , denoted as

 $[q_x, q_y, q_z] = \operatorname{IGM}(\mathbf{P}_E), \quad \mathbf{P}_E \in SE(3).$ 

## Crank & endeffector points

The crank points  $(\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_3, \mathbf{c}_4, \mathbf{c}_5, \mathbf{c}_6)$  are allowed to move on the circles defined by the motion of three actuators. The end effector points  $(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3, \mathbf{e}_4, \mathbf{e}_5, \mathbf{e}_6)$  lie on a sphere of radius *d* and center **e**.





he point parametrizations (CPL & I	EPL) are:
$\mathbf{c}_1 = [0, r \cos(q_x), l + r \sin(q_x)]^T$	$\mathbf{e}_1 = \mathbf{e} + d \cdot \mathbf{n}$

- $\mathbf{c}_2 = [0, r \cos(q_x), l r \sin(q_x)]'$
- $\mathbf{c}_3 = [l + r \sin(q_v), 0, r \cos(q_v)]^T$
- $\mathbf{c}_4 = [I r \sin(q_v), 0, r \cos(q_v)]^T$
- $\mathbf{c}_5 = [r \cos(q_z), l + r \sin(q_z), 0]^T$
- $\mathbf{e}_5 = \mathbf{e} + d \cdot \mathbf{a}$  $\mathbf{c}_6 = [r \cos(q_z), I - r \sin(q_z), 0]^T$  $\mathbf{e}_6 = \mathbf{e} - d \cdot \mathbf{a}$ .

### **Constraint equations**

xpansion of contraint equations $\ \mathbf{e}_i - \mathbf{c}_i\  =$	/ yield
$(e_x + d \cdot n_x)^2 + (e_y + d \cdot n_y - r \cdot \cos q_x)^2 + (e_z + d \cdot n_z - l - r \cdot \sin q_x)^2 = l^2$	(1)
$(e_x - d \cdot n_x)^2 + (e_y - d \cdot n_y + r \cdot \cos q_x)^2 + (e_z - d \cdot n_z - l + r \cdot \sin q_x)^2 = l^2$	(2)
$(e_x + d \cdot a_x - l - r \cdot \sin q_y)^2 + (e_y + d \cdot a_y)^2 + (e_z + d \cdot a_z - r \cdot \cos q_y)^2 = l^2$	(3)

$$(e_x - d \cdot a_x - l + r \cdot \sin q_y)^2 + (e_y - d \cdot a_y)^2 + (e_z - d \cdot a_z + r \cdot \cos q_y)^2 = l^2$$

$$(4)$$

$$(e_x + d \cdot s_x - r \cdot \cos q_z)^2 + (e_y + d \cdot s_y - l) = l^2$$

$$(5)$$

$$\frac{(e_x - d \cdot s_x + r \cdot \cos q_z)^2 + (e_y - d \cdot s_y - l}{+ r \cdot \sin q_z)^2 + (e_z - d \cdot s_z)^2 = l^2}$$
(6)

### Three virtual leg equations

By subtracting (2) from (1), (4) from (3), (6) from (5), three virtual leg equations are derived  $re_{v}\cos(q_{x})+r(e_{z}-l)\sin(q_{x})+d(ln_{z}-\mathbf{e}*\mathbf{n})=0$  $re_z \cos(q_v) + r(e_x - l) \sin(q_v) + d(la_x - \mathbf{e} * \mathbf{a}) = 0$  $re_x \cos(q_z) + r(e_y - l) \sin(q_z) + d(ls_y - \mathbf{e} * \mathbf{s}) = 0$ . With leg index  $j \in \{1, 2, 3\}$ , they are of the form (7)

 $E_i \cdot \cos(q_i) + F_i \cdot \sin(q_i) + G_i = 0$ .

### **IGM Solution**

By tangent half angle substitution  $t_j = tan(q_j/2)$ , cos  $q_j = d_j$  $(1 - t_j^2)/(1 + t_j^2)$ , sin  $q_j = 2t_j/(1 + t_j^2)$ , the equation  $(G_j - E_j) \cdot t_j^2 + 2 \cdot F_j \cdot t_j + (G_j + E_j) = 0$ in t is obtained. The two solutions for  $q_i$  are given by  $q_{j_{\pm}}, q_{j_{\pm}} = 2 \cdot \operatorname{atan2}(-F_j \pm H_j, G_j - E_j)$ 

with 
$$H_j = \sqrt{E_j^2 + F_j^2 - G_j^2}$$
, see [3]

### Rotative Inverse Geometric Model (RIGM)

Rotative Inverse Geometric Model (RIGM) is to find input joint angles for a desired orientation of the endeffector  $\mathbf{R}_E$ without knowledge of the end-effector position as

 $[q_x, q_y, q_z] = \operatorname{RIGM}(\mathbf{R}_E), \quad \mathbf{R}_E \in SO(3).$ 

 $\mathbf{e}_2 = \mathbf{e} - d \cdot \mathbf{n}$ 

 $\mathbf{e}_3 = \mathbf{e} + d \cdot \mathbf{s}$ 

 $\mathbf{e}_4 = \mathbf{e} - d \cdot \mathbf{s}$ 

Equations (1) - (6) are highly coupled. A novel iterative algorithm has been developed which can be explained by the concept of virtual joints. The method TFGM implements a three-sphere intersection to solve for e [3].

# Algorithm 1 Rotative Inverse Geometric Model (RIGM)

2:	
3:	
4:	
5:	
6:	
7:	
8:	
9:	
10:	
11:	

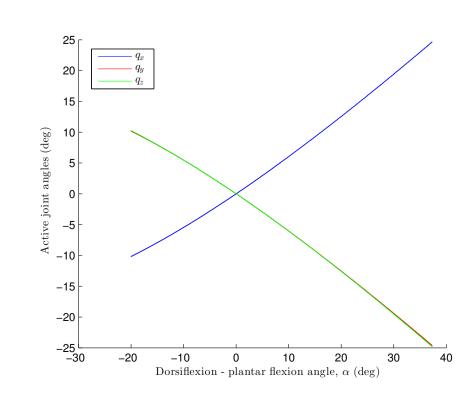
12:

During most activities of daily living, only partial ranges of motion required [4], e.g,  $10^\circ - 15^\circ$  plantar flexion and  $10^\circ$ dorsiflexion for walking on even surfaces, walking upstairs  $(37^{\circ} \text{ total ROM})$ , walking downstairs (56° total ROM).

Fig.1: Comparison of ROM between human and ACTIVE ANKLE.

# **Experimental results**

RIGM has been implemented for a task space control of the ACTIVE ANKLE. For  $\epsilon = 1.e^{-06}$  mm, the algorithm can be realized at a control frequency of 10 kHz [3].



⊳ Error

⊳ Update

### **RIGM** solution

(in) Desired orientation of the end effector,  $\mathbf{R}_E$ (out) Input joint angles  $[q_x, q_y, q_z]$  and EE shift  $[e_x, e_y, e_z]$ function RIGM( $\mathbf{R}_{E}, \epsilon$ )  $\mathbf{R}_E \mathbf{0}_{3 \times 1}$  $\mathbf{P}_E \leftarrow$ ▷ Initialization

- $\mathbf{0}_{1 \times 3}$ while  $\bar{E}_{LS} < \epsilon$  do  $(\tilde{\mathbf{e}}_1 \dots \tilde{\mathbf{e}}_6) \leftarrow \text{EPL}(\tilde{\mathbf{P}}_E)$  $[\tilde{q}_x, \tilde{q}_y, \tilde{q}_z] \leftarrow \operatorname{IGM}(\tilde{\mathsf{P}}_E)$  $(\tilde{\mathbf{c}}_1 \dots \tilde{\mathbf{c}}_6) \leftarrow \operatorname{CPL}(\tilde{q}_x, \tilde{q}_y, \tilde{q}_z)$  $E_{LS} \leftarrow \sum_{i}^{6} (\|\mathbf{\tilde{e}}_{i} - \mathbf{\tilde{c}}_{i}\| - I)^{2}$  $\tilde{\mathbf{e}} \leftarrow \mathrm{TFGM}(\tilde{q}_x, \tilde{q}_y, \tilde{q}_z, \mathbf{R}_E)$  $\mathbf{R}_E \ \mathbf{\tilde{e}}_{3 \times 1}$  $ilde{\mathbf{P}}_E \leftarrow$  $[q_x, q_y, q_z] \leftarrow [\tilde{q}_x, \tilde{q}_y, \tilde{q}_z]$
- $[e_x, e_y, e_z] \leftarrow [\tilde{e}_x, \tilde{e}_y, \tilde{e}_z]$ return  $[q_x, q_y, q_z, e_x, e_y, e_z]$

### Range of motion (ROM) comparison

Motion type	Human Ankle			ACTIVE ANKLE		
	min.	max.	abs.	min.	max.	abs.
DF – PF						
EV - IV	$-15^{\circ}$	35°	50°	$-14.62^{\circ}$	34.84°	49.46°
AD – AB	$-30^{\circ}$	45°	<b>75</b> °	$-29.20^{\circ}$	36.96°	$66.16^{\circ}$

Fig.5: Active joint angles during the dorsiflexion – plantarflexion motion.

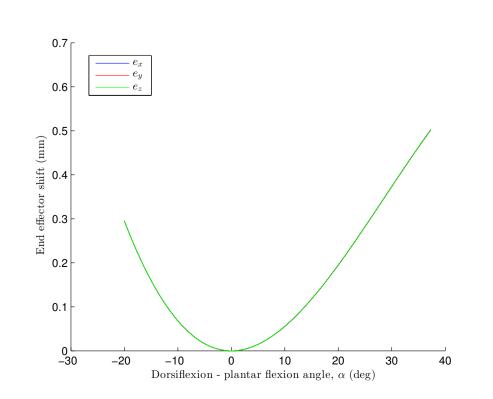


Fig.6: End effector shift during the dorsiflexion – plantarflexion motion.

### Inverse Kinematics

The three virtual leg equations (7) can be differentiated with respect to time and can be rearranged as a relation between twist  $(\mathbf{t})$  and actuated joint velocities  $(\dot{\mathbf{q}})$  through serial (**B**) and parallel (**A**) Jacobian matrices:

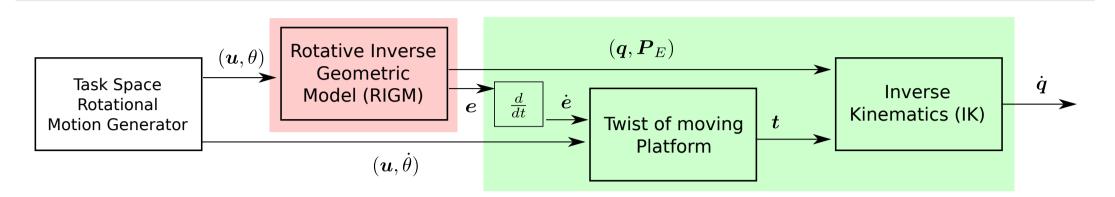


Fig.7: Cascaded task-space position and velocity control using RIGM and IK

### Task-space cascaded control

The joint-based FPGA stacks implement cascaded position, velocity and torque control. An equivalent control framework is envisioned in 3-DoF spherical task space which makes it a compact and versatile rehabilitation device.

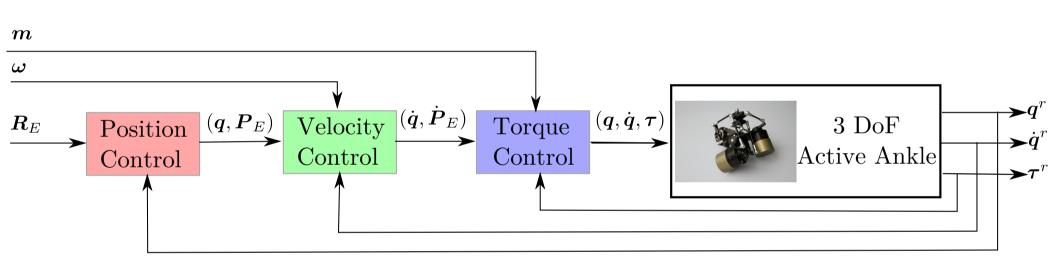


Fig.8: Cascaded task-space control scheme: a combination of desired orientation  $(\mathbf{R}_E)$ , angular velocity  $(\boldsymbol{\omega})$  and moments  $(\mathbf{m})$  in SO(3) can be the inputs.

### Conclusions

The novel ACTIVE ANKLE mechanism is briefly presented along with relevant geometric and kinematic models for its control in task-space. In the future, the cascaded task-space control framework will be equipped with torque control.

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### $\mathbf{A} \cdot \mathbf{t} = \mathbf{B} \cdot \dot{\mathbf{q}}$

The solution of the Inverse Kinematics problem requires:  $\dot{\mathbf{q}} = \mathbf{B}^{-1} \cdot \mathbf{A} \cdot \mathbf{t}$ 



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